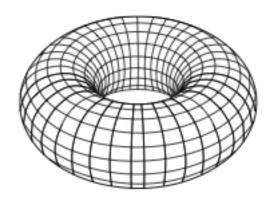
#### Manifolds & charts



A manifold **M** of dimension n is a topological space that near each point resembles n-dimensional Euclidean space  $\mathbb{R}^n$ . That is, each point of an n-dimensional manifold has a neighbourhood that is homeomorphic to the Euclidean space of dimension n

Let  $\mathcal{U} \subseteq \mathbf{M}$ , and  $\mathcal{V} \subseteq \mathbb{R}^n$ .

A homeomorphism  $\phi : \mathcal{U} \to \mathcal{V}, \ \phi(u) = (x_1(u), \dots, x_n(u))$ is a coordinate system on  $\mathcal{U}$ , and coordinates  $x_1, \dots, x_n$ . The pair  $(\mathcal{U}, \phi)$  is called a coordinate chart

Equivalence principle: choice of charts is arbitrary



• Space-time is a Lorentzian Manifold

**smooth manifold:** locally similar enough to Euclidean space to do calculus

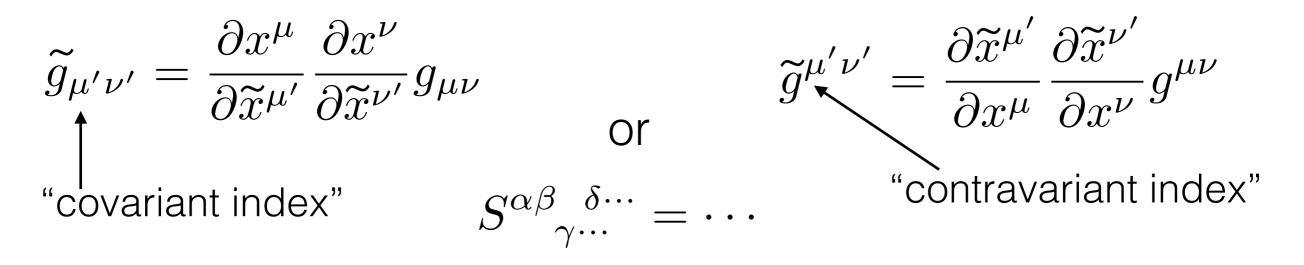
**Pseudo-Riemannian manifold:** equipped with an inner product on tangent space. Family of inner products is called Riemannian metric (tensor). Metric need not to be positive definite

Lorentzian manifold: 4 dimensional, (-,+,+,+)

tangent vectors can be classified into time-like, null, space-like

## Some basic rules

- In curved space-times we cannot use Cartesian coordinates
- Expressions have to be coordinate invariant
- coordinate transformation for tensors



•  $g^{\mu\nu}g_{\nu\sigma} = \delta^{\mu}_{\sigma}$ , i.e.,  $g^{\mu\nu}$  is the inverse of  $g_{\mu\nu}$ 



### The metric

$$\tilde{g}_{\mu'\nu'} = \frac{\partial x^{\mu}}{\partial \tilde{x}^{\mu'}} \frac{\partial x^{\nu}}{\partial \tilde{x}^{\nu'}} g_{\mu\nu}$$

- In GR the most important tensor is the metric  $g_{\mu\nu}$  which is the generalisation of  $\eta_{\mu\nu}$
- line element  $ds^2 = g_{\mu\nu} dx^{\mu} dx^{\nu}$
- there exists a single point on the manifold such that  $g_{\mu\nu} \rightarrow \eta_{\mu\nu}$  and even  $\frac{\partial}{\partial x^{\kappa}} g_{\mu\nu} \rightarrow 0$  (local inertial frame, LIF)
- $g_{\mu\nu}$  is determined by Einstein's equations
- scalar product  $g_{\mu\nu}a^{\mu}b^{\nu} = a_{\nu}b^{\nu} \equiv (a \cdot b)$



#### Derivatives

• In Euclidean space: -  $dA_i$  of vector  $A_i$  is also a vector space-independent

-  $\partial_k A_i$  form a tensor

• This is generally **not** the case in curvilinear coordinates:

- covariant vector 
$$A_{\mu} = \frac{\partial x'^{\nu}}{\partial x^{\mu}} A'_{\nu}$$
  
- its differential  $dA_{\mu} = \frac{\partial x'^{\nu}}{\partial x^{\mu}} dA'_{\nu} + A'_{\nu} \frac{\partial^2 x'^{\nu}}{\partial x^{\mu} \partial x^{\kappa}} dx^{\kappa}$ 

• Must generalise partial derivative to covariant derivative



# The covariant derivative $\nabla_{\mu}$

• Cannot be "too different" from the partial derivative. Locally free falling coordinate system (LIF):  $\nabla_{\mu} \rightarrow \partial_{\mu}$ 

$$\nabla_{\mu}V^{\nu} = \partial_{\mu}V^{\nu} + \Gamma^{\nu}_{\mu\lambda}V^{\lambda}$$

•  $\Gamma^{\nu}_{\mu\lambda}$  is the connection coefficient; result of parallel transport Locally free falling coordinate system (LIF):  $\Gamma^{\nu}_{\mu\lambda} \to 0$ 

• LIF 
$$\partial_{\mu}g_{\alpha\beta}(P) \equiv \nabla_{\mu}g_{\alpha\beta} = 0$$

Must also hold in any coordinate system!

$$\Gamma^{\alpha}_{\mu\nu} = \frac{1}{2} g^{\alpha\beta} \left( \partial_{\nu} g_{\beta\mu} + \partial_{\mu} g_{\beta\nu} - \partial_{\beta} g_{\mu\nu} \right)$$
  
Christoffel symbol  $(\Gamma^{\nu}_{\mu\lambda} = \Gamma^{\nu}_{\lambda\nu})$ 



## Covariant derivatives II

- $\Gamma^{\alpha}_{\mu\nu}$  are constructed to be non-tensorial: Instead,  $\nabla_{\mu}v^{\alpha}$  transforms as a tensor
- linearity  $\nabla(T+S) = \nabla T + \nabla S$
- product  $\nabla(T \otimes S) = (\nabla T) \otimes S + T \otimes (\nabla S)$
- On scalars we simply have  $\ 
  abla _{\mu }\Phi =\partial _{\mu }\Phi$

• On vectors 
$$\nabla_{\mu}v^{\alpha} = \partial_{\mu}v^{\alpha} + \Gamma^{\alpha}_{\mu\nu}v^{\nu}$$
$$\nabla_{\mu}v_{\alpha} = \partial_{\mu}v_{\alpha} - \Gamma^{\nu}_{\alpha\mu}v_{\nu}$$

