Problems with the Standard Model of Cosmology

Initial Conditions Dark Matter Dark Energy

The Standard Model of Cosmology



- The hot big bang + expansion of the Universe
 - explains nucleosynthesis and the CMB
- 2. + dark matter and dark energy
 - explain the growth of structures and distances to bright objects

The LCDM Model

- The standard cosmological model, LCDM, explains observations consistently in a simple framework – but we do not understand its components
 - Need *inflation* or some other theory to explain flatness of geometry and smoothness of CMB
 - We haven't detected *dark matter* and don't know what it is (it's outside the standard model of particle physics)
 - We don't know what *dark energy* is or why the value of the cosmological constant is 120 orders of magnitude off

Initial Conditions Problems

- Flatness problem: Why is $\Omega_0 = 1$? - WMAP: $\Omega_0 = 1.003 + 0.013 - 0.017$
- Horizon problem: Why is CMB so homogeneous? $-\frac{\Delta T}{T} \sim 10^{-5}$
- Density perturbations: What is the origin of $\Delta T/_T$?

The Flatness Problem

• Friedmann equation:

$$H^{2} + \frac{K}{a^{2}} = \frac{8\pi G}{3}\rho$$
$$\rightarrow |\Omega_{tot}(t) - 1| = \frac{|K|}{a^{2}H^{2}}$$

- During radiation domination: $|\Omega_{tot}(t) 1| \propto t$
- During matter domination: $|\Omega_{tot}(t) 1| \propto t^{\frac{2}{3}}$
- Thus flat geometry is *unstable* solution
 - Requires extremely small $|\Omega_{tot}(t) 1|$ in early universe

The Horizon Problem

- Due to the finite age of the Universe, the size of causally connected regions is also finite, known as the horizon
 - Co-moving horizon: $c\tau$, where τ is the conformal time
 - Physical horizon: $d_H = a(t) \int_0^t \frac{cdt'}{a(t')}$
- If $a \propto t^n$,

$$d_H = \frac{ct}{1-n} = \frac{n}{1-n} \left(\frac{c}{H}\right)$$

- Radiation dominated: a ∝ t^{1/2} → d_H = 2ct = c/H
 Matter dominated: a ∝ t^{2/3} → d_H = 3ct = 2c/H

- Today, the physical horizon is ~6 Gpc, but at the time of the last scattering surface it was much smaller
- Causally connected region of CMB subtends an angle

$$\theta \cong (1 + z_{LSS}) \left(\frac{t_{LSS}}{t_0} \right) \sim 2^{\circ}$$



Why do causally disconnected patches of CMB have the same temperature to 10^{-5} ?

• In terms of the conformal time $\tau = \int dt/a$, if $a \propto t^n$,

$$\tau = \frac{t^{1-n}}{1-n}$$



Density Perturbations

- 1. The hot big bang model does not account for origin of small initial density perturbations
- 2. There are CMB perturbations whose wavelengths are larger than the horizon at last scattering (thus acausal)



Inflation

- From the acceleration equation (neglecting Λ), $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$
- $\ddot{a} > 0$ requires w < -1/3, where $p = w\rho$
- If w = -1, $a \propto \exp(Ht)$ and

$$d_H = \frac{1}{H}(e^{Ht} - 1)$$



Some extra scalar field required to start (and stop) inflation

- <u>The Horizon Problem</u>:
- During matter and radiation domination, $\ddot{a} < 0$, so $\frac{d}{dt} \left[\frac{1}{\dot{a}} \right] = \frac{d}{dt} \left[\frac{1}{aH} \right] > 0$

Thus the Hubble radius, c/H, increases faster than the scale factor

- Objects at Hubble radius have recession velocity = c
- During acceleration, *ä* > 0, instead structures once smaller than the Hubble radius become larger and "leave the horizon"
 - No longer in causal contact; solves horizon problem



Structures enter the Hubble radius "horizon"

Structures leave the "horizon" • Acceleration also requires n > 1, where $a \propto t^n$

$$\rightarrow \tau = \int \frac{dt}{a} \propto \frac{t^{1-n}}{1+n} \rightarrow \infty$$



- How long should inflation last?
 - Physical scale of particle horizon at end of inflation must be larger than Hubble distance now:

$$\frac{1}{a_{f}H_{\inf}}e^{H_{\inf}(t_{f}-t_{i})} > \frac{1}{a_{0}H_{0}}$$

E-folding number: universe expands by factor of e^N

$$N = H_{inf}(t_f - t_i) = \ln\left(\frac{a_f}{a_i}\right) \gtrsim 60$$

– Can be very quick, e.g. $t_i = 10^{-36}s$ to $t_f = 10^{-34}s$



• <u>The Flatness Problem</u>:

$$|\Omega_{tot}(t) - 1| = \frac{|K|}{a^2 H^2}$$

- During radiation domination: $a \propto t^{\frac{1}{2}}$ $\rightarrow |\Omega_{tot}(t) - 1| \propto t$

- During inflation: $a \propto e^{Ht}$ $\rightarrow |\Omega_{tot}(t) - 1| \propto e^{-2Ht}$

• If number of e-foldings N > 60, the geometry becomes incredibly flat during inflation

- <u>Density perturbations</u>:
 - 1. Inflation gives the natural primordial perturbations: quantum fluctuations
 - 2. It is natural to have super-horizon perturbations



- <u>Initial Singularity</u>:
- t = 0 is curvature singularity: $R_{\mu\nu\rho\sigma} \rightarrow \infty$
- Energy density exceeds Planck energy,

$$E_{Pl} \equiv \sqrt{\frac{\hbar c^5}{G}} = 1.22 \times 10^{19} GeV$$

Need quantum theory of gravity

But Planck time is much earlier than inflation,

$$t_{Pl} \equiv \sqrt{\frac{\hbar G}{c^5}} = 5.39 \times 10^{-44} sec$$

Inflation independent of singularity

Dark Matter

- WIMP = weakly interacting massive particle
 - SUSY predicts as lightest super-symmetric particle
 - But SUSY disfavored by LHC?
- "WIMP miracle" correct abundance requires crosssection which is roughly what's expected for weak scale particle ~ 100 GeV



100 GeV predicts $\Omega_m \sim 0.3$



CDMS

CoGeNT

Confusing results...



• <u>Indirect detection</u>: Positron excess a hint of DM annihilation or local pulsars?



Galaxy Formation

- <u>Cusp-Core</u> <u>problem</u>: DM halo profile fits cuspy NFW in simulations but rotation curves of galaxies have a core
- Likely to be resolved by galactic physics



Galaxy Formation

- <u>Missing Satellites problem</u>: Many more satellites in CDM simulations than observed
 - Warm dark matter would suppress formation of satellites
 - Satellites may be too faint or contain only dark matter



Galaxy Formation

- <u>"Too big to fail"</u>
 <u>problem</u>: simulations
 predict many more
 massive subhalos
 than could be hosts
 of the Milky Way's
 brightest satellites
 - They aren't too faint to be observed, thus "too big to fail"



Boylan-Kolchin, et al. 2011, MNRAS, 415, L40

Dark Energy

 <u>Cosmological constant</u> <u>problem</u>: quantum field theory predicts huge value for vacuum energy, but

$$\Lambda \approx H_0^2 = (10^{-42} \,\text{GeV})^2 = 10^{-120} M_{pl}^2$$

• <u>Coincidence problem</u>: why is the era of accelerated expansion happening *now?*



Models of Acceleration

• The two main types of attempts to model acceleration depend on which side of Einstein's equation you change:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R[+?] = 8\pi G T_{\mu\nu}[-\Lambda g_{\mu\nu}?]$$

- Dark energy models add a new component such as the vacuum energy
- Modified gravity models tweak the left hand side
- More complicated models exist, or you can ignore the problem completely via the *anthropic landscape*

Models of Acceleration

• Quintessence: new degree of freedom, scalar field ϕ with potential $V(\phi)$, makes vacuum energy effectively dynamical.

$$w_q = \frac{\frac{\dot{\phi}^2}{2} - V(\phi)}{\frac{\dot{\phi}^2}{2} + V(\phi)}$$

• Many models have tracker solution: density tracks radiation until "turns on" at matter-radiation equality.



Frieman, et al. 2008, ARAA, 46, 385

Models of Acceleration

- Modified gravity: tweak GR at cosmological scales, where it has not been tested
 - Many (not all) require screening mechanism to satisfy local constraints
- Ex: f(R) gravity introduces function of Ricci scalar in GR action
 - Chameleon screening increases mass of scalar field in high density environments





The Anthropic Landscape

- The anthropic principle is an observation that the Universe we observe is consistent with the ability to produce us as observers
- Some take that as a *prediction*: various fundamental constants, etc. including the value of Λ *must* be such as to allow intelligent life
- String theory has 10⁵⁰⁰ solutions to their equations, so there is a *multiverse* and we live in the universe capable of producing us
- These arguments have *serious philosophical issues*

